

# A Preliminary Study of Space-time Finite Element Eddy-current Analysis

Junichi Niimi<sup>1</sup>, Takeshi Mifune, Tetsuji Matsuo, *Member, IEEE*

<sup>1</sup>Graduate school of Engineering, Kyoto University, niimi.junichi.57e@st.kyoto-u.ac.jp

A space-time finite element method is developed for the eddy-current analysis that can handle temporally-variant spatial finite elements arising in the analysis of moving objects. The vector potentials allocated on space-time finite elements naturally expresses the speed electromotive force. A Space-time triangular prism finite element is examined to represent the eddy-current field in space-time successfully.

*Index Terms*—Eddy current, space-time finite element, speed electromotive force

## I. INTRODUCTION

The time is usually formulated independently of the space in the finite element electromagnetic field analysis. To obtain the time evolution of electromagnetic field, a sequential scheme is generally used in a step-by-step manner with a uniform time-step. However, it is possible to handle the space and time in a unified manner in the Maxwell equations. This means that the computational finite-elements can be generated in the space-time to handle temporally-variant spatial finite elements arising in the analysis of moving objects.

Previous studies [1]-[3] have successfully introduced a temporal convolution for the symmetric formulation of space-time finite-element (FE) eddy-current analysis. However, the temporal convolution was only applied to the temporally uniform space-time elements and its application to the analysis of moving objects is not straightforward.

This paper proposes a Galerkin type space-time FE analysis to develop general polygon-type finite elements such as space-time parallelepiped and prism elements, where the speed electromotive force is naturally introduced.

## II. SPACE-TIME ALLOCATION OF VECTOR POTENTIAL

For example, the vector potential  $\mathbf{A}$  is allocated on the edges of parallelogram space-time element as in Fig. 1, where the two edges are slanted as  $(v\Delta t, \Delta t)$ . The

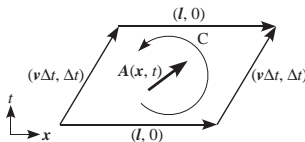


Fig.1 parallelogram space-time element

contour integral of  $\mathbf{A}$  along the parallelogram is given as

$$\begin{aligned} \oint_C \mathbf{A} \cdot d\mathbf{s} &= (v\Delta t) \cdot [\mathbf{A}(\mathbf{x} + \frac{\mathbf{l}}{2}, t) - \mathbf{A}(\mathbf{x} - \frac{\mathbf{l}}{2}, t)] \\ &\quad - \mathbf{l} \cdot [\mathbf{A}(\mathbf{x} + \frac{v\Delta t}{2}, t + \frac{\Delta t}{2}) - \mathbf{A}(\mathbf{x} - \frac{v\Delta t}{2}, t - \frac{\Delta t}{2})] \\ &\approx \mathbf{v} \cdot (\mathbf{l} \cdot \nabla) \mathbf{A} \Delta t - \mathbf{l} \cdot [\frac{\partial \mathbf{A}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{A}] \Delta t \\ &= [\mathbf{v} \cdot (\mathbf{l} \cdot \nabla) - \mathbf{l} \cdot (\mathbf{v} \cdot \nabla)] \mathbf{A} \Delta t - \Delta t \mathbf{l} \cdot \frac{\partial \mathbf{A}}{\partial t} \\ &= [(\mathbf{l} \times \mathbf{v}) \cdot \text{curl} \mathbf{A} - \mathbf{l} \cdot \frac{\partial \mathbf{A}}{\partial t}] \Delta t \\ &= \mathbf{l} \cdot [\mathbf{v} \times \text{curl} \mathbf{A} - \frac{\partial \mathbf{A}}{\partial t}] \Delta t \\ &= \mathbf{l} \cdot (\mathbf{E} + \mathbf{v} \times \mathbf{B}) \Delta t \end{aligned} \quad (1)$$

Thus, the speed electromotive force is naturally introduced, where the  $\phi=0$  gauge is used for simplicity. In the 4D space-time, the electromagnetic four-potential  $(\mathbf{A}, \phi/c)$  is defined. The contour integral of four-potential is similarly yields the speed electromotive force.

## III. SPACE-TIME FINITE ELEMENTS

Space-time finite elements can be constructed in a similar way to conventional spatial finite-elements. For example, the triangular prism element shown in Fig. 2 is examined in this study. The scalar interpolation function  $N_1$  for the node 1 given as

$$N_1 = \frac{1}{2V}(a_1 + b_1 x + c_1 t)(y_2 - y) \quad (2)$$

where  $V$  is the volume of element;  $a_1$ ,  $b_1$  and  $c_1$  are

$$a_1 = x_2 t_3 - x_3 t_2 \quad b_1 = t_2 - t_3 \quad c_1 = x_3 - x_2 \quad (3)$$

where  $x_i$  and  $t_i$  are coordinate of node  $i$ .

## IV. FORMULATION OF SPACE-TIME FINITE ELEMENT METHOD

This article examines  $xyt$ -3D space-time eddy-current

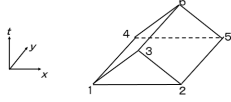


Fig.2 Triangular prism element

field described as

$$\frac{\partial}{\partial x}(\nu \frac{\partial A}{\partial x}) + \frac{\partial}{\partial y}(\nu \frac{\partial A}{\partial y}) = -J_0 + \sigma \frac{\partial A}{\partial t} \quad (4)$$

where  $A$  is the  $z$ -component of vector potential,  $J_0$  is the imposed current density and  $\nu$  is the reluctivity. The Galerkin finite-element method (FEM) is used to derive the weak form. The interpolation function  $N_i$  is multiplied to (4) as the weighting function as

$$\begin{aligned} & \iiint_{\Omega} (\frac{\partial}{\partial x}(\nu \frac{\partial A}{\partial x})N_i + \frac{\partial}{\partial y}(\nu \frac{\partial A}{\partial y})N_i) dx dy dt \\ &= \iiint_{\Omega} \sigma \frac{\partial A}{\partial t} N_i dx dy dt - \iiint_{\Omega} J_0 N_i dx dy dt . \end{aligned} \quad (5)$$

Integrating the left hand side by parts, (6) is obtained.

$$\begin{aligned} \text{(left hand side)} &= \iint \nu N_i (\frac{\partial A}{\partial x} dy - \frac{\partial A}{\partial y} dx) dt \\ &- \iiint_{\Omega} (\nu \frac{\partial A}{\partial x} \frac{\partial N_i}{\partial x} + \nu \frac{\partial A}{\partial y} \frac{\partial N_i}{\partial y}) dx dy dt \end{aligned} \quad (6)$$

The natural boundary condition is assumed to eliminate this term. Thus, the weak form (7) is obtained.

$$\begin{aligned} & \iiint_{\Omega} (\nu \frac{\partial A}{\partial x} \frac{\partial N_i}{\partial x} + \nu \frac{\partial A}{\partial y} \frac{\partial N_i}{\partial y} + \sigma \frac{\partial A}{\partial t} N_i) dx dy dt \\ &= \iiint_{\Omega} J N_i dx dy dt \end{aligned} \quad (7)$$

The linear system of equations (8) is obtained from (7)

$$([K] + [C])\{A\} = \{F\} \quad (8)$$

where

$$K_{ij} = \iiint_{\Omega} (\nu \frac{\partial N_i}{\partial x} \frac{\partial N_j}{\partial x} + \nu \frac{\partial N_i}{\partial y} \frac{\partial N_j}{\partial y}) dx dy dt \quad (9)$$

$$C_{ij} = \iiint_{\Omega} (\sigma N_i \frac{\partial N_j}{\partial t}) dx dy dt \quad (10)$$

$$F_i = \iiint_{\Omega} (J N_i) dx dy dt . \quad (11)$$

## V. ANALYSIS RESULT OF EDDY-CURRENT FIELD

The iron-cored inductor shown in Fig. 3 is analyzed by the 3D space-time FEM, where the conductivity and relative permeability of iron core are  $10^6$  S/m and 5000, and the excitation frequency is 50 Hz. The natural boundary

condition is imposed at the symmetric boundaries, and  $A = 0$  at the outer boundaries.

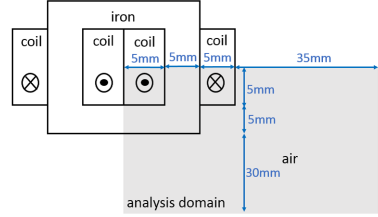


Fig.3 Iron-cored inductor

Fig. 4 shows magnetic flux lines obtained by the conventional FEM and by the space-time FEM using the triangular prism element. Both results agree. The space-time FE analysis of moving objects will be discussed in the full paper.

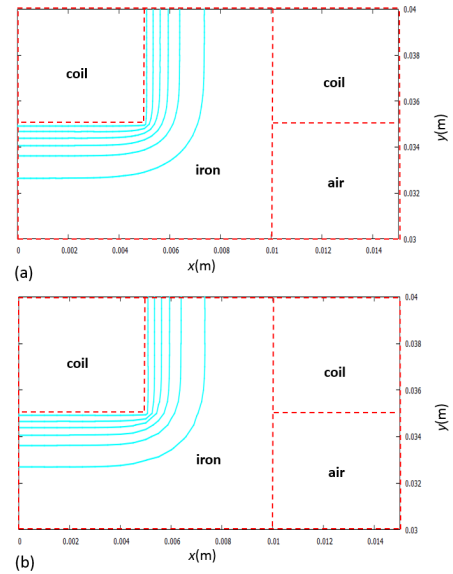


Fig.4 magnetic flux lines: (a) obtained by the conventional FEM, and (b) obtained by the space-time FEM using the triangular prism element

## VI. REFERENCES

- [1] A.J. Butler and Z.J. Cendes, "Space-time finite elements derived by convolution for the efficient solution of transient eddy current problems," IEEE Trans. Magn., vol. 24, pp. 2688-2690, Nov. 1988.
- [2] T. Renyuan, L. Feng, L. Yan and C. Xing, "Analysis of transient non-Linear eddy current fields by space-time finite element method," IEEE Trans. Magn., vol. 34, pp. 2577-2580, Sept. 1998.
- [3] S. Gyimóthy, A. Vágvölgyi, and I. Sebestyén, "Application of optimally distorted finite elements for field calculation problems of electromagnetism," IEEE Trans. Magn., vol. 38, pp. 365-368, March 2002.